THE INTERIOR ERROR OF VAN CITTERT DECONVOLUTION OF DIFFERENTIAL FILTERS IS OPTIMAL

WILLIAM LAYTON DEPARTMENT OF MATHEMATICS UNIVERSITY OF PITTSBURGH PITTSBURGH, PA 15260, USA WJL@PITT.EDU, HTTP://WWW.MATH.PITT.EDU/~WJL

Abstract. We reconsider the error in van Cittert deconvolution. We show that without any extra boundary conditions on higher derivatives of u, away from the boundary the error in van Cittert deconvolution of differential filters attains the high order of accuracy seen in the periodic problem. This error result is important for differential filters and approximate deconvolution models of turbulence.

This is an expanded version, containing more detail, background and supplementary material, of a report with the same title.

Key words. van Cittert deconvolution, singular perturbation, differential filters

1. Introduction. We consider the error in van Cittert deconvolution. We show that without any extra boundary conditions on higher derivatives of u, away from the boundary the error in van Cittert deconvolution attains the high order of accuracy seen in the periodic problem. The filtering problem is: given a function u(x) defined on a domain Ω , compute an approximation $\overline{u}(x)$ to u(x) which faithfully represents the behavior of u on scales above some, user selected, filter length (here denoted ε), and which truncates scales smaller then $O(\varepsilon)$. The deconvolution or de-filtering problem is: given \overline{u} find an accurate reconstruction of u. When the filter is smoothing $G: L^2(\Omega) \to L^2(\Omega)$ by $u \to \overline{u}$, G is compact and the deconvolution problem is illposed. One early method of deconvolution is the 1934 van Cittert [vC31] algorithm:

ALGORITHM 1.1 (van Cittert approximate deconvolution). Set $u_0 = \overline{u}$. Fix N (moderate). For n = 1, 2, ..., N - 1. perform

$$u_{n+1} = u_n + \{\overline{u} - Gu_n\}$$

Define $D_N \overline{u} := u_N$.

REMARK 1.2. The Nth van Cittert approximate deconvolution operator D_N is defined by N steps of Picard/ first order Richardson iteration for solving the operator equation $Gu = \overline{u}$ involving a possibly non-invertible operator G, [BB98]:

> given \overline{u} solve $Gu = \overline{u}$ for u. by N steps of: $u_0 = \overline{u}$ and $u_{new} = u_{old} + \{\overline{u} - Gu_{old}\}$

Since the deconvolution problem is ill-posed, convergence as $N \to \infty$ cannot expected. The relevant question is convergence for fixed N as $\varepsilon \to 0$.

The N^{th} van Cittert deconvolution operator D_N is given explicitly by

$$D_N \phi := \sum_{n=0}^{N} (I - G)^n \phi.$$
 (1.1)

The van Cittert approximate deconvolution operator corresponding to N = 0, 1, 2 and their formal orders of accuracy withy the differential filter (1.1) below are:

$$D_{0}\overline{u} = \overline{u} = u + O(\varepsilon^{2}),$$

$$D_{1}\overline{u} = 2\overline{u} - \overline{\overline{u}} = u + O(\varepsilon^{4}),$$

$$D_{2}\overline{u} = 3\overline{u} - 3\overline{\overline{u}} + \overline{\overline{\overline{u}}} = u + O(\varepsilon^{6}).$$

1

Van Cittert deconvolution thus requires only a few steps of repeated filtering. It is thus both computationally cheap and easy to program, contributing to its popularity in various applications, such as turbulence modeling, e.g., [LR12]. For convolution filters and under periodic boundary conditions, the error in van Cittert can be analyzed precisely by Fourier methods, e.g., [BIL06], [D04], [DE06]. In other cases there are significant gaps between the improved accuracy seen in computational practice and the pessimistic estimates of its global error obtained in analysis.

The goal of this report is to close this gap somewhat. We give error estimates for the van Cittert deconvolution under nonperiodic boundary conditions and preserving the boundary conditions. To develop these local and global estimates we must focus on a specific problem. To begin, we take the filter to be a differential filter (Germano [Ger86]) specifically the extension of the Pao filter, e.g., [Po00], to a bounded domain. Let Ω be a bounded, regular, planar domain with smooth boundary and $0 < \varepsilon \leq 1$ a small parameter. Given $u \in H_0^1(\Omega)$, define $Gu = \overline{u}$ as the unique solution of the elliptic-elliptic singular perturbation problem

$$-\varepsilon^2 \Delta \overline{u} + \overline{u} = u$$
, in Ω , and $\overline{u} = 0$, on $\partial \Omega$. (1.2)

Classical theory, [L73], concludes that

$$\overline{u} \to u$$
 in $L^2(\Omega)$ as $\varepsilon \to 0$.

This can be extended to $\overline{u} \to u$ in $H_0^1(\Omega)$ as $\varepsilon \to 0$, [L07]. Further, the shift theorem implies that $\overline{u} \in H_0^1(\Omega) \cap H^3(\Omega)$. Since traces of $\Delta \overline{u}$ are thus well defined, $-\varepsilon^2 \Delta \overline{u} + \overline{u} = u$ implies

$$\overline{u} = 0$$
 and $\Delta \overline{u} = 0$ on $\partial \Omega$.

As an example of the difficulties connected with the global error, consider the case N = 0 (no deconvolution) and N = 1. The regularity theory (sharp in 1d examples [L07]) predicts no improvement in the rate of convergence in L^2 , whpose norm is denoted $|| \cdot ||$. We prove the following herein which predicts improvement from higher order deconvolution in negative Sobolev norms and away from the boundary.

THEOREM 1.3 (Global and Local Deconvolution Errors). Suppose N > 0 is fixed and for k large enough that $u \in H^k(\Omega) \cap H^1_0(\Omega)$. Then

$$||u - D_0\overline{u}|| = ||u - \overline{u}|| \le C\varepsilon^2 ||u||_{H^2(\Omega)}$$

If N = 1 we have in L^2 and H^{-2}

$$||u - D_1\overline{u}|| \le C\varepsilon^2 ||u||_{H^2(\Omega)} \quad and \quad ||u - D_1\overline{u}||_{H^{-2}(\Omega)} \le C\varepsilon^4 ||u||_{H^2(\Omega)}$$

If N = 1 and additionally $\Delta u \in H_0^1(\Omega)$

$$||u - D_1 \overline{u}|| \le C \varepsilon^4 ||u||_{H^2(\Omega)}$$

If $\Delta u \neq 0$ on $\partial \Omega$ we have

$$||u - D_N \overline{u}|| \le C\varepsilon^2 ||u||_{H^2(\Omega)}$$
$$||u - D_N \overline{u}||_{H^{-2N}(\Omega)} \le C\varepsilon^{2N+2} ||u||_{H^2(\Omega)}.$$

Let $s \geq 0$. Suppose $u \in H^{2N+2}(\Omega) \cap H^1_0(\Omega)$. Let

$$\Omega_{N+1} \subset \Omega_N \subset \cdots \subset \Omega_1 \subset \Omega_0 \subset \Omega_{-1} \equiv \Omega$$

be subdomains with smooth boundaries and for $j = N + 1, \dots, 0$ with

 Ω_j has distance $C_j \varepsilon \ln(1/\varepsilon)$ from $\partial \Omega_{j-1}$,

where $C_i = C(s, N, \Omega_i, \Omega_{i-1})$. Then there is a $C = C(N, C_i)$ such that

$$||u - D_N \overline{u}||_{L^2(\Omega_{N+1})} \le C \varepsilon^{2N+2} \left[||u||_{H^{2N+2}(\Omega_0)} + \varepsilon^s ||u|| \right]$$

1.1. The case of local averaging filters. If the filter is a local averaging filter then interior estimates of the above type hold automatically because the calkculation of $\overline{\phi}$ on Ω_j only access the values of ϕ on Ω_{j-1} . Local averaging filter are vewry important in finite difference methods. Three examples follow.

Top hat filter. The top hat filter is the un-weighted average defined over a neighborhood of a given point:

$$B_{\delta}(x) := \{ y : |x - y| < \varepsilon \}, \\ \overline{u}(x) := \frac{1}{\operatorname{vol}(B_{\varepsilon}(x))} \int_{B_{\varepsilon}(x)} u(y) dy.$$

Thus in 3d this means

$$\overline{u}(x) := \frac{1}{\frac{4}{3}\pi\varepsilon^3} \int_{|x-y|<\varepsilon} u(y) dy.$$

This can be written as a convolution by choosing $g_{\delta}(x) := \varepsilon^{-3}g(x/\varepsilon)$ where

$$g(x) = 1$$
, if $|x| < \frac{3}{4\pi}$,
 $g(x) = 0$, if $|x| \ge \frac{3}{4\pi}$.

Discrete filters. In finite difference approximations, ultimately one must filter discrete velocities defined on a finite difference mesh. On a uniform mesh in 2d with averaging radius equal to the given meshwidth, $\varepsilon = \Delta x$, using the standard finite difference compass notation the analog of the top hat filter is

$$\overline{u}(P) := \frac{u(N) + u(S) + u(E) + u(W) + u(P)}{5}$$

Weighted Compact Discrete Filter, [SAK01a]. There has developed a considerable experience with weighted discrete filters inspired by the needs of difference methods. On structured meshes, filters can be derived in 1d and extended by taking tensor products of 1d filters. Generally, the higher order the filter, the more points involved in the averaging operator and thus the greater the bandwidth on the linear system that must be solved. Compact filters are a clever idea of Stolz, Adams and Kleiser [SAK01a]; they attain higher order but only require tridiagonal solves (on structured meshes). The following 1d weighted discrete filter from [SAK01a] has second order accuracy and has proven its value in large eddy simulation. Given values u_i of the variable u at equi-spaced mesh points x_i , a weighting parameter α is chosen in the range $-1/2 \leq \alpha \leq +1/2$. Then, filtered values are calculated by solving the tridiagonal system

$$\alpha \overline{u}_{i-1} + \overline{u}_i + \alpha \overline{u}_{i+1} = \left(\frac{1}{2} + \alpha\right) \left(u_i + \frac{u_{i-1} + u_{i+1}}{2}\right).$$

The fact that the inverse of a tridiagonal matrix is a full matrix means that local error estimates for van Cittert deconvolution do ot follow automatically for weighted compact filters. This extension is an open problem.

2. Proof of the deconvolution error estimate. Since van Cittert deconvolution is mathematically equivalent to a truncation of a geometric (operator) series, it is quite easy to calculate the deconvolution error for specific choices of filter for smooth functions. The error in van Cittert deconvolution is thus calculated, [BIL06], [D04], [DE06], to be

$$u - D_N \overline{u} = (I - G)^{N+1} u = (-1)^{N+1} \varepsilon^{2N+2} \triangle^{N+1} G^{N+1} u$$

$$= O(\varepsilon^{2N+2}) \text{ for } C^{\infty}_{periodic} \text{ functions } u.$$
(2.1)

Thus, accuracy of van Cittert in any norm $||| \cdot |||$ depends on whether, and for what values on N,

$$||\triangle^{N+1}G^{N+1}(u)||| \le C(u) < \infty$$
 uniformly in ε .

The proof is based on the error representation (2.1) and two regularity results for the elliptic-elliptic singular perturbation problem. The global regularity result was proven in [L07], see also [LR12]. The local, interior regularity result is a special case of Theorem 2.3, page 26 of Navert [N82] (setting the convecting velocity to zero), see also [SW83]. We shall forst recall these two results, give a preliminary lemma and them give the proof (which is short with this preparation). $H^k(\Omega)$ denotes the Sobolev space of all functions with derivatives of order $\leq k$ in $L^2(\Omega)$. The $L^2(\Omega)$ norm is $|| \cdot ||$ and $H_0^1(\Omega) := \{v \in H^1 : v = 0 \text{ on } \partial\Omega\}$. For (1.1) we assume (in particular implying u = 0 on $\partial\Omega$)

$$u \in H^k(\Omega) \bigcap H^1_0(\Omega). \tag{2.2}$$

This condition precludes simple boundary layers in \overline{u} but does not imply higher derivatives of \overline{u} are free of layers. From (1.1) it also implies that $\Delta \overline{u} = 0$ on $\partial \Omega$.

THEOREM 2.1 (Theorem 1.1 in [L07]). Suppose $u \in H^2(\Omega) \cap H^1_0(\Omega)$. Then there is a constant C > 0 independent of ε such that

$$||\overline{u}||_{H^{1}(\Omega)} \leq C||u||_{H^{1}(\Omega)}$$
, for $l = 0, 1, 2.$ (2.3)

If $u \in H^4(\Omega) \cap H^1_0(\Omega)$, $\Delta u \in H^1_0(\Omega)$. Then

$$||\overline{u}||_{H^{l}(\Omega)} \leq C||u||_{H^{l}(\Omega)} , \text{ for } l = 0, 1, 2, 3, 4.$$

$$(2.4)$$

In general, suppose $u \in H^{2k}(\Omega) \cap H^1_0(\Omega), \Delta^j u \in H^1_0(\Omega), j = 1, \dots, k-1$. Then for $l = 1, \dots, 2k$

$$||\overline{u}||_{H^{1}(\Omega)} \leq C||u||_{H^{1}(\Omega)}.$$
(2.5)

Examples in [L07] show that the estimate $||\overline{u}||_l \leq C||u||_l$, being limited to l = 0, 1, 2 is sharp unless higher derivatives of u are zero on $\partial\Omega$.

THEOREM 2.2 (Special case of Navert [N82], Theorem 2.3). For $u \in H^k(\Omega) \cap H^1_0(\Omega)$ consider

$$-\varepsilon^2 \Delta \overline{u} + \overline{u} = u$$
, in Ω , and $\overline{u} = 0$, on $\partial \Omega$. (2.6)

Let $m \ge 0, s \ge 0$. Let $\Omega' \subset \Omega'' \subset \Omega$ be subdomains with smooth boundaries with

• Ω' has distance $C_1 \varepsilon \ln(1/\varepsilon)$ from $\partial \Omega''$,

• Ω'' has distance $C_2 \varepsilon \ln(1/\varepsilon)$ from $\partial \Omega$.

where $C_i = C_i(s, m, \Omega', \Omega'')$. Then the solution to (1.1) satisfies

$$|\overline{u}||_{H^m(\Omega')} \le C\left(||u||_{H^m(\Omega'')} + \varepsilon^s ||u||\right)$$

First we calculate the global regularity of repeated filtering. PROPOSITION 2.3. Let $u \in H^k(\Omega) \cap H^1_0(\Omega)$. We have for $J \ge 1$

$$||G^{J}u||_{H^{k}(\Omega)} \leq C||G^{J-1}u||_{H^{k}(\Omega)}, k = 0, 1, \cdots, 2J$$

Proof. For n = 1 Theorem 1.1 implies

$$||\overline{u}||_{H^k(\Omega)} \leq C||u||_{H^k(\Omega)}$$
, for $k = 0, 1, 2$ and $\Delta \overline{u} = 0$ on $\partial \Omega$.

Since $\Delta \overline{u} = 0$ on $\partial \Omega$, we repeat. Indeed, $G^2 u = G \overline{u} = \overline{\overline{u}}$ so that

$$\begin{aligned} ||\overline{\overline{u}}||_{H^{k}(\Omega)} &\leq C ||\overline{u}||_{H^{k}(\Omega)} \text{, for } k = 0, 1, 2, 3, 4 \text{ and that} \\ \triangle \overline{\overline{u}} &= \triangle \overline{u} = 0 \text{ on } \partial \Omega. \end{aligned}$$

Taking the Laplacian of the equation for $\overline{\overline{u}}$ gives

$$-\delta^2 \triangle^2 \overline{\overline{u}} + \triangle \overline{\overline{u}} = \triangle \overline{u} , \text{ in } \Omega.$$
(2.7)

Now, let $x \to \partial \Omega$ and use $\Delta \overline{\overline{u}} = \Delta \overline{\overline{u}} = 0$ on $\partial \Omega$. This implies

$$\triangle^2 \overline{\overline{u}} = \triangle \overline{\overline{u}} = \overline{\overline{u}} = 0 \text{ on } \partial \Omega$$

so that for $\overline{\overline{u}}$.

$$||\overline{\overline{u}}||_{H^{k}(\Omega)} \leq C||\overline{\overline{u}}||_{H^{k}(\Omega)}$$
, for $k = 0, 1, 2, 3, 4, 5, 6$.

The proof continues by induction. \Box

We can now prove the deconvolution error estimate in Theorem 1.1.

Proof. [Proof of Theorem 1.1] We consider $\triangle^{N+1}G^{N+1}(u)$ and use Theorem 1.1 in [L07] repeatedly. For N = 0 this is $||\triangle(-\varepsilon^2 \triangle + 1)^{-1}u||$:

$$||u - D_0\overline{u}|| = ||u - \overline{u}|| = \varepsilon^2 ||\triangle(-\varepsilon^2 \triangle + 1)^{-1}u||.$$

The first estimate follows since:

$$|\triangle(-\varepsilon^2 \triangle + 1)^{-1}u|| = ||\triangle \overline{u}|| \le C||\overline{u}||_2 \le C||u||_2$$

For N = 1 and under $\Delta u \in H_0^1(\Omega)$ we have similarly that $||\overline{\overline{u}}||_4 \leq C||u||_4$. Thus

$$||u - D_1\overline{u}|| = \varepsilon^4 ||\Delta^2\overline{\overline{u}}|| \le C\varepsilon^4 ||\overline{\overline{u}}||_4 \le C\varepsilon^4 ||u||_4$$

For the H^{-2} estimate we use that $\triangle^2 \overline{\overline{u}} = \triangle (\overline{\Delta \overline{u}})$. Step by step, using $\triangle \overline{u} = 0$ on $\partial \Omega$ we find $|| \triangle^2 \overline{\overline{u}} ||_{-2} \leq C || \overline{\Delta \overline{u}} || \leq C || \Delta \overline{u} || \leq C || u ||_2$, completing the proof. The case of N > 1 follows the same way.

For the interior estimates we use Theorem 2.2 as follows.

$$||u - D_N \overline{u}||_{L^2(\Omega_{N+1})} = \varepsilon^{2N+2} ||(\Delta^{N+1} G^{N+1})u||_{L^2(\Omega_{N+1})} \le \le C \varepsilon^{2N+2} ||G^{N+1}u||_{H^{2N+2}(\Omega_{N+1})}.$$

Note that $||\overline{\phi}|| \leq ||\phi||$ so that $||G^j u|| \leq ||u||$ for all j. Now $G^{N+1}u = \overline{\phi}, \phi = G^N u$. Thus, for any s > 0

$$\begin{aligned} ||G^{N+1}u||_{H^{2N+2}(\Omega_{N+1})} &\leq C\left(||G^{N}u||_{H^{2N+2}(\Omega_{N})} + \varepsilon^{s}||G^{N}u||\right) \\ &\leq C\left(||G^{N}u||_{H^{2N+2}(\Omega_{N})} + \varepsilon^{s}||u||\right). \end{aligned}$$

We repeat this argument. Indeed, $G^N u = \overline{\phi}, \phi = G^{N-1}u$. Thus, for any s > 0

$$||G^{N}u||_{H^{2N+2}(\Omega_{N})} \le C\left(||G^{N-1}u||_{H^{2N+2}(\Omega_{N-1})} + \varepsilon^{s}||u||\right)$$

At the last step we have, for any s > 0

$$||G^{1}u||_{H^{2N+2}(\Omega_{1})} \leq C\left(||u||_{H^{2N+2}(\Omega_{0})} + \varepsilon^{s}||u||\right).$$

Thus (recalling that N is fixed and C can depend on N) we have

$$||u - D_N \overline{u}||_{L^2(\Omega_{N+1})} \le C\varepsilon^{2N+2} \left[||u||_{H^{2N+2}(\Omega_0)} + \varepsilon^s ||u|| \right]$$

г		
	_	

3. Conclusions. The filtering or convolution operator $G: u \to \overline{u}$ is a bounded map: $L^2(\Omega) \to L^2(\Omega)$. If (as for differential filtyers) it is smoothing, its inverse cannot be bounded due to small divisor problems. Indeed, it is known quite generally that inversion is not well posed.

THEOREM 3.1. Let H be a Hilbert space and $G: H \to H$ a compact map. Then, if H is infinite dimensional

$$Range(G) \neq H.$$

In other words, G is not invertible as a bounded linear operator.

Thus stable exact deconvolution is not possible and approximate deconvolution must be used instead. An approximate deconvolution operator D_N is an approximate inverse $\overline{u} \to D_N(\overline{u}) \approx u$ which

- is a bounded operator on $L^2(\Omega)$,
- approximates u in some useful (typically asymptotic) sense, and
- satisfies other conditions necessary for the application at hand.

The error in van Cittert approximate deconvolution in the non periodic case is of high accuracy, away from boundaries, like that of the periodic case. It is an interesting analytic open question to establish if a similar result holds for the Stokes differential filter. It is also an interesting algorithmic open question to alter the van Cittert procedure near boundaries to obtain a high order accurate reconstruction of the unknown function up to the boundary.

REFERENCES

[BIL06]	L. C. BERSELLI, T. ILIESCU, AND W. LAYTON, Mathematics of Large Eddy Simulation of Turbulant Flows, Springer Barlin, 2006
[BJG07]	L.C. BERSELLI, V. JOHN AND C. GRISANTI, Analysis of commutation errors for func- tions with low regularity I Comput Appl Math 206 (2007) 1027-1045
[BL11]	L.C. BERSELLI AND R. LEWANDOWSKI, Convergence of approximate deconvolution models to the filtered Navier-Stokes Equations to appear in Ann IHP 2011
[B79]	J. BARANGER, On the thickness of the boundary layer in elliptic-elliptic singular perturbation problems, 395-400 in: Numerical Analysis of Singular Perturbation Problems (PW Hemker and LLH Miller, eds.) Academic press NY, 1979
[B75]	J.G. BESJES, Singular perturbation problems for linear elliptic differential operators of arbitrary order L Degeneration to elliptic operators JMAA 40(19795) 24-46
[D04]	A. DUNCA, Space averaged Navier-Stokes equations in the presence of walls, Ph.D. Thesis, University of Pittsburgh, 2004.
[DE06]	A. DUNCA AND Y. EPSHTEYN, On the Stolz-Adams deconvolution model for the Large- Eddy simulation of turbulent flows, SIAM J. Math. Anal., 37(6) (2006), 1890- 1902.
[E79]	W. ECKHAUS, Asymptotic analysis of singular perturbations, N. Holland, Amster- dam, 1979.
[GT]	D. GILBARG AND N.S. TRUDINGER, Elliptic partial differential equations of second order, Springer, Berlin, 2001.
[Ger86] [Geu97]	 M. GERMANO, Differential filters of elliptic type, Phys. Fluids, 29(1986), 1757-1758. B. J. GEURTS, Inverse modeling for large eddy simulation, Phys. Fluids, 9 (1997), 3585.
[ILT05]	A. A. ILYIN, E. M. LUNASIN AND E. S. TITI, A modified Leray-alpha subgrid-scale model of turbulence, Nonlinearity, 19 (2006), 879-897.
[L07]	W. LAYTON, A remark on regularity of elliptic-elliptic singular perturbation prob- lem, Technical Report, available at http://www.math.pitt.edu/techreports.html, (2007).
[LL03]	W. LAYTON AND R. LEWANDOWSKI, A simple and stable scale similarity model for large eddy simulation: energy balance and existence of weak solutions, Applied Math. letters 16(2003) 1205-1209
[LL05]	W. LAYTON AND R. LEWANDOWSKI, Residual stress of approximate deconvolution large eddy simulation models of turbulence. Journal of Turbulence, 46(2): 1-21, 2006
[LL06a]	W. LAYTON AND R. LEWANDOWSKI, On a well posed turbulence model, Discrete and Continuous Dynamical Systems - Series B, 6(2006) 111-128.
[LMNR06b]	W. LAYTON, C. MANICA, M. NEDA AND L. REBHOLZ, Numerical analysis of a high ac- curacy Leray-deconvolution model of turbulence, to appear: Numerical Methods for PDEs, 2007.
[LMNR06]	W. LAYTON, C. MANICA, M. NEDA AND L. REBHOLZ, The joint Helicity- Energy cascade for homogeneous, isotropic turbulence generated by approx- imate deconvolution models, University of Pittsburgh Technical Report, http://www.math.pitt.edu/techreports.html, 2006.
[LN06a]	W. LAYTON AND M. NEDA, Truncation of scales by time relaxation, Journal of Mathematical Analysis and Applications, 325(2): 788-807, 2007.
[LN06b]	W. LAYTON AND M. NEDA, The energy cascade for homogeneous, isotropic turbulence generated by approximate deconvolution models, to appear in JMAA, 2007.
[LR12]	W. LAYTON AND L. REBHOLZ, Approximate deconvolution models of turbulence. Springer LNM, Springer, Berlin, 2012.
[Le06]	R. LEWANDOWSKI, Vorticities in a LES model for 3D periodic turbulent flows, Journ. Math Fluid Mech Vol 8 pp 398-422 2006
[LP09]	R. LEWANDOWSKI AND Y. PREAUX, Attractors for a deconvolution model of turbulence, Applied Mathematics Letters 22 (2009) 642-645
[L73]	JL. LIONS, Perturbations singulieres dans les problemes aux limites et en controle antimal Springer LNM vol 323–1973
[MR10]	 W. MILES AND L. REBHOLZ, An enhanced physics based scheme for the NS-alpha turbulence model, Numerical Methods for Partial Differential Equations, 26(6), (2010) 1530-1555
[N82]	U. NÄVERT, A finite element method for convection diffusion problems, Thesis, Chalmers Inst. of Technology, 1982.

[N10]	M. NEDA, Discontinuous time relaxation for the time dependent Navier-Stokes equa- tions. Advances in Numerical Analysis, 2010 (ID 419021) (2010), 1-21
[Po00]	S POPE Turbulent Flows Cambridge Univ Press (2000)
[Reb07]	L. REBHOLZ, Conservation laws of turbulence models, Journal of Mathematical Analysis and Applications, 326(1) (2007), 33-45.
[S01] [SW83]	 P. SAGAUT, Large eddy simulation for Incompressible flows, Springer, Berlin, (2001). A SCHATZ AND L WAHLBIN, On the finite elementy method for singularly perturbed reaction- diffusion problems in two and one spoace domensions, Math. Comp. 40(1983), 47-89.
[SA99]	S. STOLZ AND N. A. ADAMS, On the approximate deconvolution procedure for LES, Phys. Fluids, 11 (1999), 1699-1701.
[SAK01a]	S. STOLZ, N. A. ADAMS AND L. KLEISER, The approximate deconvolution model for LES of compressible flows and its application to shock-turbulent-boundary-layer interaction, Phys. Fluids 13 (2001), 2985-3001.
[SAK01b]	S. STOLZ, N. A. ADAMS AND L. KLEISER, An approximate deconvolution model for large eddy simulation with application to wall-bounded flows, Phys. Fluids, 13 (2001), 997-1015.
[SAK02]	S. STOLZ, N. A. ADAMS AND L. KLEISER, The approximate deconvolution model for compressible flows: isotropic turbulence and shock-boundary-layer interaction, in: Advances in LES of complex flows (editors: R. Friedrich and W. Rodi) Kluwer, Dordrecht, (2002).
[SSK05]	S. STOLZ, P. SCHLATTER, AND L. KLEISER, <i>High-pass filtered eddy-viscosity models</i> for LES of transitional and turbulent flow, Phys. Fluids, 17 (2005)065103.
[St03]	B. STRAUGHAN, The Energy Method, Stability and Nonlinear Convection, Springer, Berlin, (2003).
[T93]	L. TARTAR, Remarks on some interpolation spaces, in: BVPs for PDEs and applica- tions, Masson, Paris, 1993, pp. 229-252.
[vC31]	P. VAN CITTERT, Zum Einfluss der Spaltbreite auf die Intensitats verteilung in Spek- trallinien II, Zeit. fur Physik 69 (1931), 298-308.

3.1. More References to work on approximate deconvolution models in CFD.

REFERENCES

[AS01]	N. A. ADAMS AND S. STOLZ, Deconvolution methods for subgrid-scale approximation in large eddy simulation, Modern Simulation Strategies for Turbulent Flow, R.T.
[AS02]	 Edwards, 2001. N. A. ADAMS AND S. STOLZ, A subgrid-scale deconvolution approach for shock cap- turing Journal of Computational Physics 178 (2002) 391-426
[BGJ07]	L.C. BERSELI, C.R. GRISANTI, AND V. JOHN, Analysis of commutation errors for functions with low regularity, J. Comput. Appl. Math. 206 (2007), 1027-1045.
[BB98]	M. BERTERO AND B. BOCCACCI, Introduction to Inverse Problems in Imaging, IOP Publishing Ltd., (1998).
[BR10]	A. BOWERS AND L. REBHOLZ, Increasing accuracy and efficiency in FE computations of the Leray-deconvolution model, Numerical Methods for Partial Differential Equations, to appear, (2010).
[BLCR10]	A. BOWERS, B. COUSINS, A. LINKE AND L. REBHOLZ, New connections between finite element formulations of the Navier-Stokes equations, Journal of Computational Physics 229 (2010) 9020-9025
[Cag10]	A. CAGLAR, Convergence analysis of the Navier-Stokes-alpha model, Numerical meth- ods for partial differential equations, 26 (2010), 1154-1167.
[CHOT05]	A. CHESKIDOV, D. D. HOLM, E. OLSON AND E. S. TITI, On a Leray-α model of turbulence, Royal Society London, Proceedings, Series A, Mathematical, Physical and Engineering Sciences, 461 (2005), 629-649.
[CKG01]	S. CHILDRESS, R. R. KERSWELL AND A. D. GILBERT, Bounds on dissipation for Navier-Stokes flows with Kolmogorov forcing, Phys. D., 158 (2001), 1-4.
[Con10]	J. CONNORS, Convergence analysis and computational testing of a finite element discretization of the Navier-Stokes alpha model, Numerical Methods for Partial Differential Equations, 26(6) (2010), 1328-1350.
[DM01]	A. DAS AND R.D. MOSER, Filtering boundary conditions for LES and embedded boundary simulations, DNS/LES progress and challenges (C. Liu, L. Sakeland,
	8

[D03]	and T. Beutner, eds.), Greyden Press, Columbus, 2001, pp. 389{396. A. DUNCA, Optimal design of fluid flow using subproblems reduced by large eddy simulation Technical Report ANL/MCS (2003)
[ELN06]	V. ERVIN, W. LAYTON AND M. NEDA, Numerical analysis of a higher order time relaxation model of fluide Int. I. Numer, Anal. and Modeling, 4(3-4) (2007)
[FHT01]	648-670. C. FOIAS, D. D. HOLM AND E. S. TITI, The Navier-Stokes-alpha model of fluid
[FHT09]	turbulence, Physica D, (152-153) (2001), 505-519.
[11102]	tions, and their relation to the Navier-Stokes equations and turbulence theory, Journal of Dynamics and Differential Equations, 14 (2002), 1-35.
[GL00]	G. P. GALDI AND W. J. LAYTON, Approximation of the large eddies in fluid motion II: A model for space-filtered flow, Math. Models and Methods in the Appl. Sciences,
[G03]	 10 (2000), 343-350. B. J. GEURTS, <i>Elements of direct and large eddy simulation</i>, Edwards Publishing, (2003)
[GH05]	B. J. GEURTS AND D. D. HOLM, Leray and LANS-alpha modeling of turbulent mixing, L. of Turbulence, 00(2005), 1-42
[GH03]	B. J. GEURTS AND D. D. HOLM, Regularization modeling for large eddy simulation, Physics of Fluids, 15(1) (2003) 13-16
[Gue04]	 R. GUENANFF, Non-stationary coupling of Navier-Stokes/Euler for the generation and radiation of aerodynamic noises, Ph.D. thesis, Dept. of Mathematics, Uni-
[Guer]	versite Rennes 1, Rennes, France, (2004). JL. GUERMOND, Subgrid stabilization of Galerkin approximations of monotone op- erators, C. R. Acad. Sci. Paris, S érie I, 328(7) (1999), 617-622.
[GOP03]	J.L. GUERMOND, S. PRUDHOMME AND J.T. ODEN, An interpretation of the NS alpha model as a frame indifferent Leray regularization, Physica D Nonlinear Phenom-
[HV03]	ena, 177 (2003), 23-30. HASELBACHER, A. AND VASILYEV, O.V., Commutative discrete filtering on unstruc- tured grids based on least-squares techniques, Journal of Computational Physics,
[HAD06]	 187(1) (2003), 197-211. S. HICKEL, N.A. ADAMS, J.A. DOMARADZKI, An adaptive local deconvolution method for implicit LES 1, Comput. Phys. 213 (2006)413-436
[J04]	V. JOHN, Large Eddy Simulation of Turbulent Incompressible Flows, Springer, Berlin, (2004)
[JL01]	V. JOHN AND W. LAYTON, Approximating local averages of fluid velocities, Computing 66 (2001) 269-287.
[LLMN08]	A. LABOVSCHII, W. LAYTON, C. MANICA, M. NEDA, L. REBHOLZ, I. STANCULESCU AND C. TRENCHEA, Architecture of Approximate Deconvolution Models of Turbulence, Ercoftac Series: Quality and Reliability of Large-Eddy Simulation, 10.1007/978- 1-4020-8578-9_1, Editors: Johan Meyers, Bernard J. Geurts and Pierre Sagaut, (2008).
[L10]	W LAYTON, Existence of smooth attractors for the Navier-Stokes-omega model of turbulence, JMAA, 366, (2010), 81-89.
[L07b]	W. LAYTON, Superconvergence of finite element discretization of time relaxation mod- els of advection, BIT Numerical Mathematics 47 (2007), 565-576.
[LL05]	W. LAYTON AND R. LEWANDOWSKI, Residual stress of approximate deconvolution large eddy simulation models of turbulence. Journal of Turbulence, 46(2) (2006), 1-21
[LMNR09]	 W. LAYTON, C. MANICA, M. NEDA, AND L. REBHOLZ, Numerical analysis and compu- tational comparisons of the NS-alpha and NS-omega regularizations, Computer Mathods in Applied Machanics and Engineering, 199 (2010), 916-931
[LRS10]	 W. LAYTON, L. REBHOLZ AND M. SUSSMAN, Energy and helicity dissipation rates of the NS-alpha and NS-alpha-deconvolution models, IMA Journal of Applied Mathematics 75(1) (2010) 56-74
[LS07]	W. LAYTON AND I. STANCULESCU, K-41 optimized approximate deconvolution models, International Journal of Computing Science and Mathematics 1 (2007) 206 411
[LS09]	 W. LAYTON AND I. STANCULESCU, Chebychev optimized approximate deconvolution models of turbulence, Applied Mathematics and Computation, 208 (2009), 106- 112
[LST08]	W. LAYTON, I. STANCULESCU, AND C. TRENCHEA, Theory of the NS-\overline{\overlin}\overlin{\overline{\overline{\overline{\overline{\overlin}\overlin{\overline{\overlin}\overlin{\verline{\overlin{\

[LLe06]	J. LEDERER AND R. LEWANDOWSKI, On the RANS 3D model with unbounded eddy viscosities Ann IHP ann non lin 24(3) (2007) 413-441
[MM06]	 C. C. MANICA AND S. KAYA-MERDAN, Convergence Analysis of the Finite Element Method for a Fundamental Model in Turbulence, University of Pittsburgh Tech- nical Report. http://www.math.pitt.edu/techreports.html (2006)
[MNOR11]	C.C. MANICA, M. NEDA, M.A. OLSHANSKII, L. REBHOLZ AND N. WILSON, On an efficient finite element method for NS- $\overline{\omega}$ with strong mass conservation, Computational Methods in Applied Mathematics, to appear, (2011).
[MS09]	C. C. MANICA AND I. STANCULESCU, Leray-Tikhonov regularization mod- els of fluid motion, University of Pittsburgh Technical Report, http://www.math.pitt.edu/techreports.html, (2009).
[PTGG06]	C. D. PRUETT, B. C. THOMAS, C. E. GROSCH, AND T. B. GATSKI, A temporal approximate deconvolution model for large-eddy simulation, Phys. Fluids, 18 (2006), 1-4.
[RS10]	L. REBHOLZ AND M. SUSSMAN, On the high accuracy NS-alpha-deconvolution model of turbulent fluid flow, Mathematical Models and Methods in Applied Sciences, 20(4) (2010), 611-633.
[S08]	I. STANCULESCU, Existence theory of abstract approximate deconvolution models of turbulence, Ann. Univ. Ferrara, 54 (2008), 145-168.
[VLM98]	O. VASILYEV, T. LUND AND P. MOIN, A general class of commutative filters for LES in complex geometries, Journal of Computational Physics, 146 (1998), 105-123.
[VTC05]	M. I. VISHIK, E. S. TITI AND V. V. CHEPYZHOV, Trajectory attractor approximations of the 3d Navier-Stokes system by the Leray-alpha model, Russian Math Dokladi, 71 (2005), 91-95.
[V03]	A.W. VREMAN, The filtering analog of the variational multiscale method in large-eddy simulation, Phys Fluids 15(2003) 61-64.